

Cite this: *Nanoscale*, 2015, 7, 12133

## Fabrication of SnO<sub>2</sub>–SnO nanocomposites with p–n heterojunctions for the low-temperature sensing of NO<sub>2</sub> gas†

Lei Li,<sup>a,b</sup> Chunmei Zhang<sup>a,c</sup> and Wei Chen<sup>\*a</sup>

In this report, the fabrication of a novel SnO<sub>2</sub>–SnO nanostructure with p–n heterojunctions has been achieved through a facile one-pot and low-cost hydrothermal process. The structure and properties of the nanocomposite were analyzed with X-ray techniques and electron microscopy. HRTEM characterization showed that the p–n heterojunctions were formed with small n-type SnO<sub>2</sub> nanocrystals dispersed on the surface of large p-type SnO crystals. Compared to the single SnO<sub>2</sub>-based material, a gas sensor fabricated from the SnO<sub>2</sub>–SnO composite exhibited an enhanced sensing performance for NO<sub>2</sub> gas detection, with a limit of detection and sensitivity of 0.1 ppm and 0.26 ppm<sup>-1</sup>, respectively, at a relatively low operating temperature (50 °C). Moreover, the p–n heterojunctions exhibited high sensing selectivity for NO<sub>2</sub>. Such a high sensing sensitivity and a low operating temperature make the SnO<sub>2</sub>–SnO p–n nanomaterial a promising gas sensor for practical NO<sub>2</sub> gas detection. The improved sensing response characteristics of the hybrid material could be attributed to the p–n junctions formed through the *in situ* growth of SnO<sub>2</sub> nanocrystals on SnO nanoplates. The present study is helpful for the design of novel gas sensing materials and the development of NO<sub>2</sub> gas sensors.

Received 12th April 2015,  
Accepted 3rd June 2015  
DOI: 10.1039/c5nr02334c

www.rsc.org/nanoscale

### 1. Introduction

Tin dioxide (SnO<sub>2</sub>) nanomaterials have been widely applied in the field of gas sensing for decades due to their unique sensing properties, such as high sensitivity, fast response and high stability, *etc.*<sup>1,2</sup> Tin dioxide is a typical n-type semiconductor material with rutile crystalline structure and a wide band gap (3.6 eV). Many endeavors have been devoted to the design of SnO<sub>2</sub> nanostructures, and have resulted in many excellent achievements.<sup>1,2</sup> In particular, owing to their excellent tunable physicochemical properties and their size-, shape- and exposed surface-dependent gas sensing properties, SnO<sub>2</sub> nanomaterials with various structures have been designed for gas sensing, such as zero-dimensional (0D) particles<sup>3,4</sup> or clusters,<sup>5</sup> one-dimensional (1D) nanofibers<sup>6,7</sup> or nanorods,<sup>8,9</sup> two-dimensional (2D) nanosheets<sup>10</sup> and three-dimensional (3D) hierarchical and porous hollow structures,<sup>11–14</sup> *etc.* In order to

overcome the intrinsic shortcomings of pure SnO<sub>2</sub> nanomaterials and to improve the sensing performance, many SnO<sub>2</sub> nanomaterials doped with different dopants have been proposed recently. Among the doped materials, the semiconductor materials doped or loaded with noble metals such as Pt,<sup>15–18</sup> Pd,<sup>19–22</sup> Au<sup>23–25</sup> and Ag<sup>26,27</sup> have been synthesized, and the doped SnO<sub>2</sub> materials exhibited enhanced sensing performances. However, due to the high cost and the easy poisoning of noble metals by toxic gases,<sup>22,28,29</sup> the application of noble metal-doped SnO<sub>2</sub> nanomaterials in gas sensors is largely limited. In recent years, it was found that hybrid sensing materials fabricated from different types of semiconductors by forming p–n heterojunctions with SnO<sub>2</sub> nanomaterials can effectively increase charge transfer and enhance the sensing properties.<sup>30–33</sup>

The unique heterojunction structure fabricated from SnO<sub>2</sub> and other doped or supported metal oxide materials with various morphologies has exhibited outstanding gas sensing properties. Compared with gas sensors based on a single SnO<sub>2</sub> material, nano-gas sensors fabricated from p–n heterojunction structures of SnO<sub>2</sub>/metal oxides can largely overcome the shortcomings of high resistance, high working temperatures and long response times, *etc.*<sup>30,33</sup> In a p–n junction structure, between the contact sections of the n- and p-type semiconductor materials, the conduction and valence bands usually bend and the Fermi levels are equalized associated with the for-

<sup>a</sup>State Key Laboratory of Electroanalytical Chemistry, Changchun Institute of Applied Chemistry, Chinese Academy of Sciences, Changchun 130022 Jilin, People's Republic of China. E-mail: weichen@ciac.ac.cn; Tel: +86 431-85262723

<sup>b</sup>The College of Electrical and Electronic Engineering, Changchun University of Technology, Changchun 130012, People's Republic of China

<sup>c</sup>University of Chinese Academy of Sciences, Beijing 100039, People's Republic of China

† Electronic supplementary information (ESI) available: Survey XPS spectrum of the SnO<sub>2</sub>–SnO hybrid material. See DOI: 10.1039/c5nr02334c

mation of depletion layers, which directly results in the improvement of the conductivity and the acceleration of response and recovery times.<sup>34,35</sup> Moreover, the morphologies of metal oxides with hierarchical structures can be tuned to enhance the gas diffusion rate and enlarge the specific surface area, and thus to further improve the gas sensing properties. In recent years, typical SnO<sub>2</sub>-based nanomaterials with p-n heterojunctions have been designed as gas sensors, e.g. SnO<sub>2</sub> nanoparticle-doped WO<sub>3</sub> nanolamellae for NO<sub>2</sub> sensing,<sup>36</sup> α-Fe<sub>2</sub>O<sub>3</sub> nanorods grown epitaxially on the surface of SnO<sub>2</sub> nanosheets for acetone sensing,<sup>37</sup> heterostructured CuO nanoparticles and SnO<sub>2</sub> nanowires for H<sub>2</sub>S detection,<sup>34</sup> ZnO/SnO<sub>2</sub> core-shell structural microspheres for ethanol sensing,<sup>38</sup> and so on. Meanwhile, heterostructures of SnO<sub>2</sub> metal oxide materials are not only confined to gas sensing applications, but also applied in electrochemical sensors,<sup>39</sup> photocatalysts,<sup>40</sup> lithium ion batteries,<sup>41</sup> dye-sensitized solar cells,<sup>42</sup> etc.

Nitrogen dioxide (NO<sub>2</sub>) is a reddish-brown, high-activity gaseous substance, which is mainly released from high-temperature combustion processes, such as the exhaust gas emissions of automobiles, exhaust gas discharged from power stations and so on.<sup>43,44</sup> It is worth mentioning that NO<sub>2</sub> is also a key component of acid rain, which causes a variety of environmental effects, including visibility reduction in the atmosphere, surface water acidification and increased harmful toxins in fish and other aquatic organisms in water. Therefore, efficient sensing materials for the sensitive detection of NO<sub>2</sub> gas in the ambient with a fast response are essentially required.<sup>45</sup> In this paper, to develop practically useful NO<sub>2</sub> sensing materials, SnO<sub>2</sub>-SnO hybrids with p-n junctions were fabricated through a one-step hydrothermal process. To the best of our knowledge, there is no previous report on the synthesis and NO<sub>2</sub> gas sensing properties of SnO<sub>2</sub>-SnO hybrid materials with p-n junctions. The prepared SnO nanoplate-SnO<sub>2</sub> nanoparticle composites can effectively enhance the electron transfer and improve the response to the target NO<sub>2</sub> gas. The gas sensing tests showed that detection of NO<sub>2</sub> with a limit of detection of 0.1 ppm can be realized at a low temperature (50 °C) by using the prepared SnO<sub>2</sub>-SnO p-n junctions as sensing probes. The present study demonstrates that designing novel semiconductor-semiconductor nanostructures with p-n heterojunctions is an efficient route to develop next-generation gas sensors with improved sensing performance.

## 2. Experimental section

### 2.1 Materials preparation

All reagents were of analytical grade and were used without further purification unless otherwise specified. The SnO<sub>2</sub>-SnO nanocomposites were synthesized by a hydrothermal process. Typically, 0.066 g of SnCl<sub>2</sub>·2H<sub>2</sub>O (Beijing Chemical Reagent Co., Ltd, China, AR, 98%) and 0.054 g of urea were dissolved in 15 mL of deionized water under moderate stirring for about 30 min until the mixture appeared as a homogeneous milky solution. The mixed solution was then transferred to a 25 mL

Teflon-lined stainless steel vessel and the hydrothermal process was performed at 180 °C for 16 h. The formed SnO<sub>2</sub>-SnO composites were washed with water and ethanol by centrifugation (3000 rpm), and then dried at 40 °C for several hours in air to obtain the final product. In order to compare the gas sensing performances of the prepared SnO<sub>2</sub>-SnO nanocomposites and single SnO<sub>2</sub> material, SnO<sub>2</sub> nanoparticles were also prepared by the same synthetic method and treatment procedure. In a typical synthesis, 0.105 g of SnCl<sub>4</sub>·5H<sub>2</sub>O (Chengdu Gracia Chemical Technology Co., Ltd, China, AR, 99%), and 0.054 g of urea were dispersed in 15 mL deionized water under magnetic stirring. Notably, this mixed solution was colorless and clear, which was different from that of SnCl<sub>2</sub>·2H<sub>2</sub>O. The subsequent hydrothermal and purification processes were the same as that aforementioned.

### 2.2 Materials characterization

The composition and crystalline structures of the products were characterized by X-ray powder diffraction (XRD) on a Bruker D8 Advance X-ray diffractometer between 10° and 90° using a Cu Kα radiation source (λ = 1.54 Å). The morphologies and selected-area electron diffraction of the composites were characterized by transmission electron microscopy (TEM, JEOL 2000) and high resolution TEM (HRTEM, JEM-2010(HR) microscope). At the same time, the chemical states of the products were examined by X-ray photoelectron spectroscopy (XPS, VG Thermo ESCALAB 250 spectrometer). The Brunauer-Emmett-Teller (BET, Quantachrome Autosorb Automated Gas Sorption System) specific surface area and porosity properties analysis were also performed. UV-vis spectra were collected on a UV-3000PC Spectrophotometer (Shanghai Mapada Instruments Co., Ltd). Finally, the gas sensing properties of the materials were tested using a CGS-1TP (Chemical Gas Sensor-1 Temperature Pressure) intelligent gas sensing analysis system (Beijing Elite Tech Co., Ltd, China).

### 2.3 Gas sensing measurements

The gas sensing behavior of the samples was tested on a heating ceramic plate of the CGS-1TP instrument by pressing the sample firmly with two probes. The products were first mixed with *N,N*-dimethylformamide (DMF) to form homogeneous suspensions, which were then drip-dropped on commercial ceramic substrates with Ag-Pd interdigitated electrodes. The as-prepared samples were dried at 55 °C during the process of drip-dropping in air. The gas sensing measurements were conducted with the prepared sensors at different operating temperatures with a fixed working voltage of 9 V. The operating temperature was changed by an external temperature controller under the heating ceramic plate and could be adjusted from room temperature to 500 °C with a rate of 3–5 °C s<sup>-1</sup> (the relative humidity was about 38%). The sensor was preheated for about 30 min at a certain target temperature until the resistance of the materials reached a stable range, and it was then sealed in an 18 L test chamber. The target gas was injected into the test chamber through the injection pore. During the measurements, the fans for gas

mixing were turned on until the resistance of the material reached a stable range. After the gas sensing tests, the test chamber was opened to expose the sensor to fresh air and the fans were turned off. These procedures were repeated for the gas sensing tests under different conditions. Here, the response of the sensor is defined as  $R = R_g/R_a$ , where  $R_g$  and  $R_a$  are the resistances of the sensor in the target gas and in air, respectively. The concentration of the target gas could be calculated using eqn (1):<sup>43</sup>

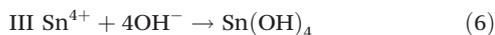
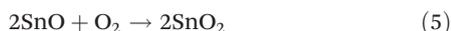
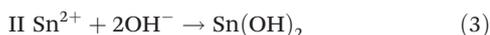
$$Q = V \times C \times 10^{-6} \times \frac{273 + T_R}{273 + T_B}, \quad (1)$$

where  $Q$  is the volume of the target gas,  $V$  is the volume of the test chamber,  $C$  is the concentration of the target gas,  $T_B$  is the operating temperature and  $T_R$  is the room temperature.

### 3. Results and discussion

#### 3.1 Synthesis and characterization of the gas sensing materials

In this study, two kinds of NO<sub>2</sub> gas sensing materials (SnO<sub>2</sub>-SnO nanocomposites and single SnO<sub>2</sub> nanoparticles) were prepared by one-pot hydrothermal processes. In the hydrothermal process, urea plays a key role for the formation of the SnO<sub>2</sub>-SnO composites. When the temperature is higher than 90 °C, urea begins to hydrolyze and NH<sub>4</sub><sup>+</sup> and CO<sub>3</sub><sup>2-</sup> are slowly released. With further increase in temperature, CO<sub>2</sub> and OH<sup>-</sup> can be produced from CO<sub>3</sub><sup>2-</sup> and NH<sub>4</sub><sup>+</sup>, respectively.<sup>46</sup> Owing to the limited oxygen amount in the solution, only a small part of the SnO can be converted to SnO<sub>2</sub> to form the SnO<sub>2</sub>-SnO composite. While totally oxygen-free conditions are hard to achieve in the hydrothermal process, pure SnO product cannot be prepared with the present method. The main reactions from the Sn<sup>2+</sup> and Sn<sup>4+</sup> precursors could be described as follows:



The compositions and crystalline structures of the obtained products were first characterized by X-ray techniques and electron microscopy. As shown in Fig. 1A, the XRD pattern of the heterojunction-structured SnO<sub>2</sub>-SnO composite agrees well with standard cassiterite (SnO<sub>2</sub>, JCPDS card no. 41-1445) and romarchite (SnO, JCPDS card no. 06-0395). However, the product synthesized from SnCl<sub>4</sub>·5H<sub>2</sub>O shows only a cassiterite SnO<sub>2</sub> phase (JCPDS card no. 41-1445). It is noteworthy that, compared with the diffraction peaks of the pure SnO<sub>2</sub> sample, the intensity and full widths at half-maximum (FWHM) of the diffraction peaks from the SnO<sub>2</sub>-SnO composite are much

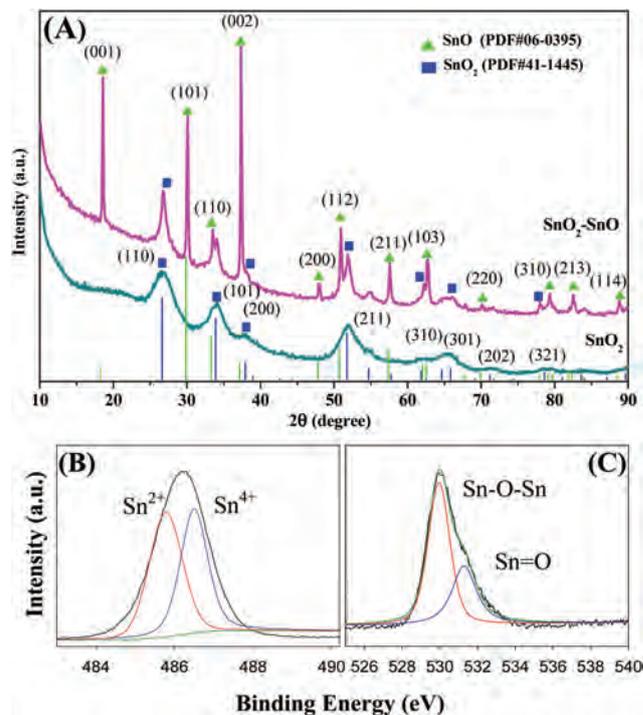


Fig. 1 (A) Typical XRD patterns of the SnO<sub>2</sub>-SnO nanocomposite (purple) and SnO<sub>2</sub> nanoparticles (turquoise). For comparison, the XRD data of bulk SnO (green bars) and SnO<sub>2</sub> (blue bars) from the joint Committee Powder Diffraction Standard are also included. High resolution XPS spectra of Sn 3d<sub>5/2</sub> (B) and O 1s (C) of the SnO<sub>2</sub>-SnO nanocomposite.

stronger and narrower, respectively. These results imply that the SnO<sub>2</sub> nanocrystals formed in the SnO<sub>2</sub>-SnO composite have higher degrees of crystallinity and relatively larger crystal sizes than the pure SnO<sub>2</sub> sample. On the other hand, such results suggest that during the formation of SnO<sub>2</sub>-SnO hybrid materials, nucleation and crystal growth can be promoted. Interestingly, in the XRD pattern of SnO<sub>2</sub>-SnO, the intensity and FWHM of the diffraction peaks from SnO are much stronger and narrower than those from the SnO<sub>2</sub> component, suggesting that the SnO nanocrystals are much larger than the SnO<sub>2</sub> nanoparticles. Meanwhile, from the standard XRD data of SnO (JCPDS card no. 06-0395), the (101) diffraction peak has the highest intensity ( $I_{(001)}/I_{(101)} = 0.1$ ,  $I_{(002)}/I_{(101)} = 0.14$ ). However, in the XRD pattern of the SnO<sub>2</sub>-SnO composite, the (001) and (002) diffraction peaks are much stronger than the (101) peak, with the much increased  $I_{(001)}/I_{(101)}$ ,  $I_{(002)}/I_{(101)}$  ratios of 1.1 and 1.26. Such XRD features strongly suggest that the formed SnO nanocrystals in SnO<sub>2</sub>-SnO are predominated by an exposed (001) facet. The lattice parameters of the (001), (101) and (002) facets were calculated to be  $c = 0.485$ ,  $0.485$  and  $0.436$  nm, respectively, based on the MDI Jade 5.0 software. The crystallite sizes of SnO and SnO<sub>2</sub> in the SnO<sub>2</sub>-SnO composite, and that of the pure SnO<sub>2</sub> nanoparticles, were evaluated to be 80.64, 28.4 and 4 nm, respectively, by the Scherrer equation:  $d = 0.9\lambda/B_{2\theta} \cos\theta$  ( $\lambda$ : wavelength of the

X-rays;  $\theta$ : angle of the peak;  $B_{2\theta}$ : the peak width at half-height), with the values of the FWHM of the major diffraction peaks of the SnO (002) and SnO<sub>2</sub> (110) planes. The calculated crystalline sizes are in accordance with the following results from the TEM and HRTEM measurements.

XPS measurements were also carried out to identify the chemical states of the SnO<sub>2</sub>-SnO composite. Fig. S1† shows the full scale survey spectrum of the sample. Clearly, the peaks from Sn, O and C can be observed without the presence of any other impurities. From the high-resolution Sn 3d<sub>5/2</sub> spectrum of the SnO<sub>2</sub>-SnO composite shown in Fig. 1B, the two fitted peaks at 485.8 and 486.5 eV correspond to the Sn<sup>2+</sup> and Sn<sup>4+</sup> states, respectively.<sup>30,47</sup> Fig. 1C shows the high-resolution O 1s spectrum of the hybrid, in which the deconvoluted two peaks at binding energies of 529.9 and 531.2 eV correspond to the SnO<sub>2</sub> (Sn-O-Sn) and to the oxygen of SnO (Sn=O) bonds, respectively.<sup>48</sup> The XPS results are in good agreement with those of the XRD measurements. Both XRD and XPS characterization indicate that the SnO<sub>2</sub>-SnO hybrid has been successfully fabricated through the hydrothermal process.

Fig. 2A and B show TEM images of the as-prepared pure SnO<sub>2</sub> nanoparticles at different magnifications. It can be seen that the obtained SnO<sub>2</sub> nanoparticles have tiny sizes. Due to

the absence of any surfactants and supporting materials, the produced SnO<sub>2</sub> nanoparticles exhibit aggregates with a porous structure. Meanwhile, the SnO<sub>2</sub>-SnO hybrids obtained by using different tin precursors through the same hydrothermal process exhibit a very different morphology from that of the pure SnO<sub>2</sub> material. Fig. 2C and D show TEM images of the SnO<sub>2</sub>-SnO composite at different magnifications. It can be seen that in the composite, the SnO crystal exhibits a laminar structure of several hundreds and tens of nanometers in lateral size and thickness, respectively. At the same time, one layer of the SnO<sub>2</sub> nanoparticles of approximately 4–10 nm in diameter is dispersed on the surface of the large SnO plates, forming the heterojunction structures. To further examine the p-n heterojunction structure, a typical HRTEM image of the SnO<sub>2</sub>-SnO hybrid is shown in Fig. 2E. Highly crystalline SnO<sub>2</sub> nanocrystals with clear lattice fringes from different facets can be observed. Meanwhile, the lattice fringes corresponding to the (001) plane of a SnO nanoplate can also be seen. A selected-area electron diffraction (SAED) pattern is shown in Fig. 2F. It can be seen that, in addition to the diffraction rings from multiple SnO<sub>2</sub> nanocrystals, there are also apparent diffraction spots, indicating the single crystalline phase of the SnO nanoplates. These electron microscopy analyses are in accordance with the above XRD and XPS measurements, showing the formation of SnO<sub>2</sub>-SnO p-n structures with large SnO laminar crystals covered by small SnO<sub>2</sub> nanoparticles.

In order to quantitatively determine and compare the specific surface areas and pore volumes of the SnO<sub>2</sub> and SnO<sub>2</sub>-SnO hybrid, multipoint BET surface area and pore volume measurements were carried out based on the nitrogen adsorption-desorption isotherms. The N<sub>2</sub> physisorption isotherms of the two samples and their corresponding Barrett-Joyner-Halenda (BJH) pore size distribution curves are shown in Fig. 3. One can see that the SnO<sub>2</sub>-SnO hybrid exhibits an isotherm feature of the H3 type according to the IUPAC classification, whereas the pure SnO<sub>2</sub> sample shows a H2 type. From the measurements, the BET surface areas of SnO<sub>2</sub>-SnO and SnO<sub>2</sub> are 45.6 and 178.9 m<sup>2</sup> g<sup>-1</sup>, respectively. As described above, the high surface area of the pure SnO<sub>2</sub> nanocrystals could be attributed to the small particle size and the porous structure which can be seen in the TEM characterization. On the basis of the BJH model and the desorption data, the average pore radii  $dV(d)$  of the SnO<sub>2</sub>-SnO composite and SnO<sub>2</sub> nanoparticles were calculated to be 1.68 and 4.3 nm, respectively, as shown in the Fig. 3 insets.

### 3.2 Sensing performances of the SnO<sub>2</sub>-SnO hybrid and pure SnO<sub>2</sub> nanoparticles for NO<sub>2</sub> gas detection

The sensing performances of the SnO<sub>2</sub>-SnO hybrid and single SnO<sub>2</sub> for NO<sub>2</sub> gas detection were studied and compared. First, the operating temperature is one of the important sensing properties of materials, and a low operating temperature is especially favorable for their practical application in real gas sensors. Fig. 4 shows the responses of the two materials towards 50 ppm of NO<sub>2</sub> gas as a function of the operating temperature. It can be seen that for both materials, the

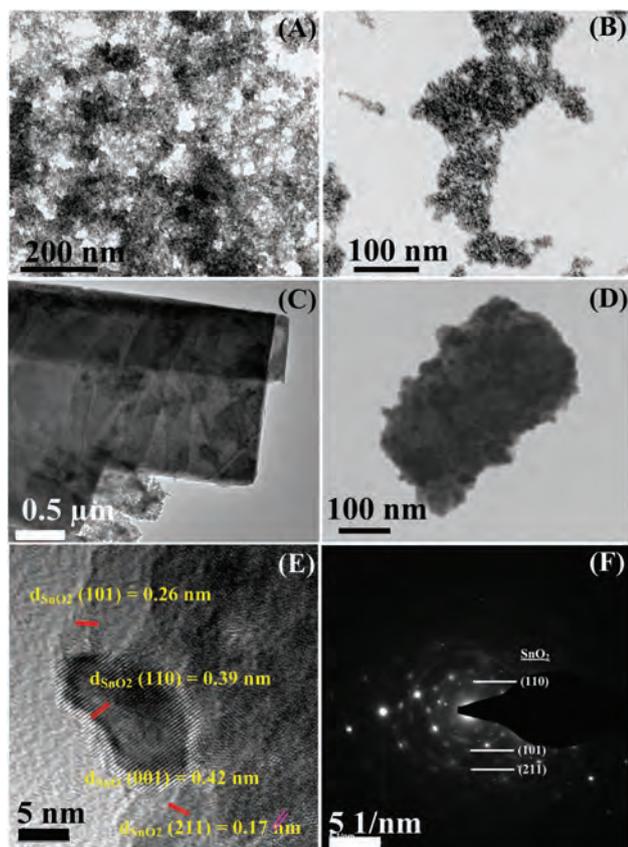


Fig. 2 (A, B) TEM images of the SnO<sub>2</sub> nanoparticles at different magnifications. (C, D) TEM images of the SnO<sub>2</sub>-SnO hybrid at different magnifications. High-resolution TEM image (E) and a selected-area electron diffraction (SAED) pattern (F) of the SnO<sub>2</sub>-SnO hybrid.

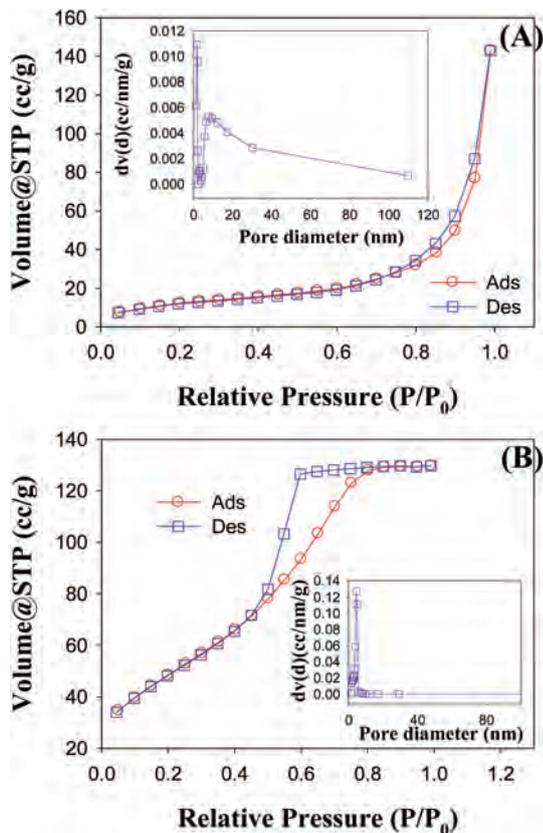


Fig. 3  $N_2$  adsorption-desorption isotherms of the  $SnO_2$ - $SnO$  composite (A) and pure  $SnO_2$  nanoparticles (B). The insets are the corresponding pore size distributions.

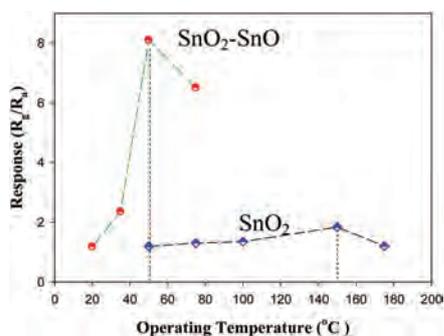


Fig. 4 Sensing responses of  $SnO_2$ - $SnO$  and  $SnO_2$  to 50 ppm  $NO_2$  gas as a function of the operating temperature.

response intensity is dependent on the working temperature, *i.e.* increasing first with the appearance of a peak, and then declining with further increasing temperature. However, the response intensity and the optimum operating temperature are very different for the two materials. For the  $SnO_2$ - $SnO$  hybrid, the sensing response gradually increases first with increasing temperature and reaches a maximum at 50 °C, and even at lower temperatures, *ca.* room temperature, the material

still shows a relatively high response. Such a result implies that after further optimization, sensors based on the  $SnO_2$ - $SnO$  hybrid could have the sensing ability for  $NO_2$  gas detection at room temperature. This is a big breakthrough for semiconductor metal oxide-based gas sensors, since  $NO_2$  gas sensing can be achieved at room temperature without any aid from noble metal doping or large surface-area supports (*e.g.* graphene, carbon nanotubes). However, the sensing response of the pure  $SnO_2$  reaches a maximum at around 150 °C, which is triple that obtained from the  $SnO_2$ - $SnO$  hybrid (50 °C). Moreover, the response intensity from pure  $SnO_2$  is much lower than that from the  $SnO_2$ - $SnO$  composite (about one fourth of that of the hybrid). From the XRD and TEM measurements, the particle size of pure  $SnO_2$  is around 4 nm and the size of the  $SnO_2$  nanoparticles on the surface of  $SnO$  is about 28 nm. According to the size effect, compared with larger nanoparticles, smaller-sized particles should have better sensing performances due to their larger surface area.<sup>49</sup> However, the present study shows that the  $SnO_2$ - $SnO$  hybrid with larger particle size exhibits a much higher sensing performance. Therefore, the enhanced sensing performance of the  $SnO_2$ - $SnO$  hybrid should be ascribed to the unique p-n heterojunction structure. Such excellent gas sensing properties with an enhanced sensing response and a much lowered operating temperature make the novel  $SnO_2$ - $SnO$  hybrid structure a kind of promising gas sensor material for  $NO_2$  gas detection.

The dynamic sensing response and recovery properties of the pure  $SnO_2$  and  $SnO_2$ - $SnO$  for  $NO_2$  detection were investigated. Fig. 5A and B show the sensing responses of  $SnO_2$  and  $SnO_2$ - $SnO$  as a function of the concentration of  $NO_2$  gas (from 15 to 110 ppm) at 150 and 50 °C, respectively. From Fig. 5A, it can be seen that with the increase in  $NO_2$  concentration, the resistance of pure  $SnO_2$  exhibits little change, indicating the non-sensitivity of pure  $SnO_2$  even at high temperature (150 °C). However, as shown in Fig. 5B, at a much lower temperature (50 °C), the resistance of the  $SnO_2$ - $SnO$  hybrid shows an increase with increasing  $NO_2$  gas concentration. Moreover, the resistance could be almost completely recovered when the sensor was exposed to fresh air during the whole testing procedure. Such results indicate the high sensing sensitivity of the  $SnO_2$ - $SnO$  hybrid toward  $NO_2$  gas at a low temperature and suggest that the sensing mechanism of the  $SnO_2$ - $SnO$  composites corresponds to that of typical n-type semiconductor metal oxides.

Interestingly, we also found that the pure  $SnO_2$  exhibited different sensing characteristics at low and high temperatures. As shown in Fig. 5C, at 75 °C, the introduction of 50 ppm of  $NO_2$  gas led to an increase in the resistance of the pure  $SnO_2$ . However, at 175 °C, the presence of  $NO_2$  gas can result in a decrease in the resistance of  $SnO_2$ , which is totally opposite to the behavior at low temperature. In other words, pure  $SnO_2$  exhibits n-type characteristics at low temperature (75 °C) but shows an overturned sensing behavior at 175 °C with a p-type sensing mechanism, which is different from those reported previously.<sup>34,45</sup> We will discuss this issue in the later section about the gas sensing mechanism. The responses of the  $SnO_2$ -

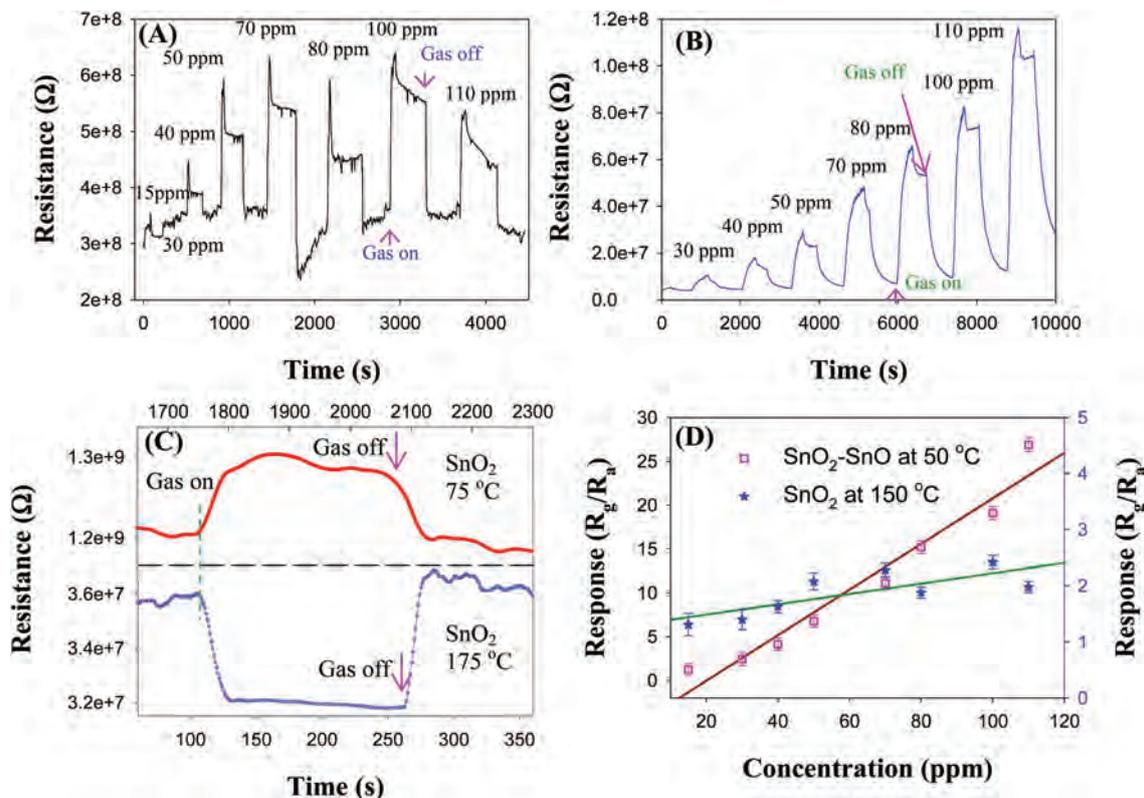


Fig. 5 Sensing responses of pure SnO<sub>2</sub> at 150 °C (A) and SnO<sub>2</sub>-SnO at 50 °C (B) upon exposure to NO<sub>2</sub> gas with different concentrations varying from 15 to 110 ppm. (C) Sensing responses of pure SnO<sub>2</sub> to 50 ppm NO<sub>2</sub> gas at 75 °C (top) and 175 °C (bottom). (D) Response variations of SnO<sub>2</sub>-SnO and SnO<sub>2</sub> as a function of NO<sub>2</sub> concentration at 50 and 150 °C, respectively. Solid lines show the linear fitting of the experimental data.

SnO hybrid and the pure SnO<sub>2</sub> as a function of NO<sub>2</sub> concentration are plotted in Fig. 5D. For the SnO<sub>2</sub>-SnO hybrid in the whole concentration range, a good linear fit was obtained with  $R^2 = 0.943$ . From the slope of the calibration line, the sensitivity of the SnO<sub>2</sub>-SnO hybrid for NO<sub>2</sub> detection was calculated to be  $0.257 \text{ ppm}^{-1}$ . Furthermore, the limit of detection (LOD) of the SnO<sub>2</sub>-SnO hybrid was estimated to be 0.1 ppm based on the relation  $\text{LOD} = 3 \alpha_D/\alpha$  (where  $\alpha$  is the slope of the linear part of the calibration curve and  $\alpha_D$  is the standard deviation of the noise in the response curve in air).<sup>50</sup> In sharp contrast, a linear fit with  $R^2 = 0.555$  was obtained from the pure SnO<sub>2</sub> nanoparticles. Compared with the high sensitivity and low LOD of the SnO<sub>2</sub>-SnO hybrid, the pure SnO<sub>2</sub> material is not sensitive to the concentration variation of NO<sub>2</sub> gas (green fitting line) with sensitivity and theoretical LOD of  $0.007 \text{ ppm}^{-1}$  and 0.6 ppm, respectively. Meanwhile, it can be seen that the SnO<sub>2</sub> sensor reaches sensing saturation at a very low concentration, and thus the concentration change of the target gas can be hardly distinguished. These results strongly indicate that the SnO<sub>2</sub>-SnO hybrid has an enhanced sensing performance for NO<sub>2</sub> gas compared to that of pure SnO<sub>2</sub> nanomaterials. Table 1 summarizes the sensing performances of SnO<sub>2</sub>-based hybrid metal oxide nanomaterials with p-n heterostructures for the detection of different gases. Clearly, in this study, the optimum temperature, limit of detection and

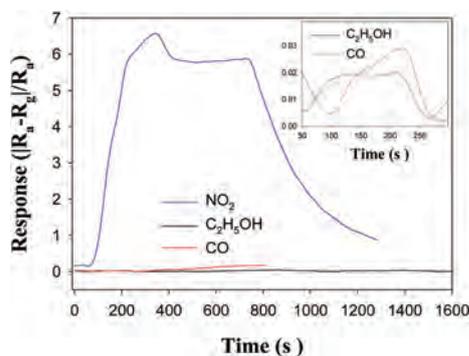
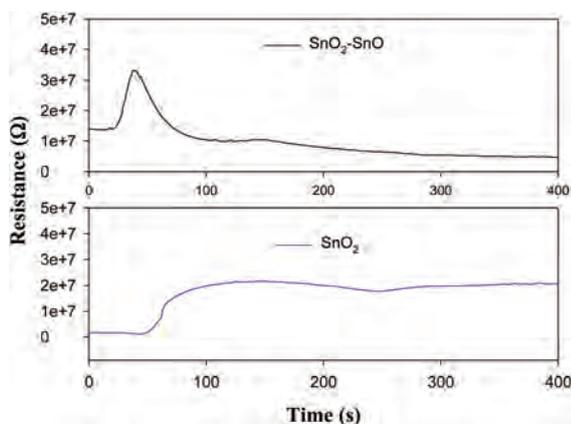
response intensity to 100 ppm target gas almost exceed those in the previous reports. The sensing selectivity of the SnO<sub>2</sub>-SnO hybrid for NO<sub>2</sub> gas was also examined. As shown in Fig. 6, with SnO<sub>2</sub>-SnO exposed to 100 ppm NO<sub>2</sub>, C<sub>2</sub>H<sub>5</sub>OH and CO at 50 °C, the response intensity to NO<sub>2</sub> was significantly larger than those to C<sub>2</sub>H<sub>5</sub>OH and CO gases, indicating the high selectivity of the p-n heterojunction nanostructure for NO<sub>2</sub> detection. The significantly enhanced sensing performance of the SnO<sub>2</sub>-SnO hybrid makes it a promising sensing material for applications in real NO<sub>2</sub> gas sensors.

### 3.3 Gas sensing mechanism of the SnO<sub>2</sub>-SnO composite with p-n heterojunctions

From the above studies, the SnO<sub>2</sub>-SnO hybrid with p-n junctions exhibited improved sensing properties compared to the pure SnO<sub>2</sub> material. Here, we propose a sensing mechanism of the SnO<sub>2</sub>-SnO p-n heterojunction nanostructure. The semiconductive properties of the two materials were first studied.<sup>51</sup> Fig. 7 shows the variation of the resistances of the hybrid SnO<sub>2</sub>-SnO and pure SnO<sub>2</sub> materials in air as a function of time in the temperature range from ~20 to 200 °C. We can see that these two sensing materials show different variation tendencies, although they belong to the same type of semiconductor metal oxides (n-type). SnO<sub>2</sub> is a well known n-type semiconductor (the majority of carriers are electrons). Thus,

**Table 1** Summary of the gas sensing performances of SnO<sub>2</sub>-based metal oxide nanomaterials with p–n heterostructures

Analyte gas	Composition	Synthesis route	Morphology	Opera. temp. (°C)	Response	Concentration	LOD	Ref.
NO <sub>2</sub>	WO <sub>3</sub> -SnO <sub>2</sub>	Combinatorial solution deposition	SnO <sub>2</sub> particles loaded on WO <sub>3</sub> lamellae	200	N/A	200 ppb	N/A	36
C <sub>2</sub> H <sub>5</sub> OH	ZnO-SnO <sub>2</sub>	Hydrothermal method	Core-shell	250	52.7	50 ppm	N/A	38
CH <sub>3</sub> COCH <sub>3</sub>	Fe <sub>2</sub> O <sub>3</sub> -SnO <sub>2</sub>	Hydrothermal method	Core-shell	370	1.7	10 ppm	10 ppm	56
CH <sub>3</sub> COCH <sub>3</sub>	Fe <sub>2</sub> O <sub>3</sub> -SnO <sub>2</sub>	Hydrothermal method	Fe <sub>2</sub> O <sub>3</sub> nanorods grown on SnO <sub>2</sub> nanosheets	250	17	100 ppm	N/A	37
H <sub>2</sub> S	CuO-SnO <sub>2</sub>	Chemical vapour deposition	CuO particles onto SnO <sub>2</sub> nanowires	200	26.3	10 ppm	N/A	34
H <sub>2</sub> S	CuO-SnO <sub>2</sub>	Electrospun	Composite nanowires	300	650	100 ppm	N/A	31
(CH <sub>3</sub> ) <sub>2</sub> CHOH	SnO <sub>2</sub> -SnO <sub>2</sub>	Chemical precipitation	Mixed phase SnO <sub>2</sub> nanorods	255	4.7	100 ppm	N/A	57
H <sub>2</sub>	SnO <sub>2</sub> -SnO	Hydrothermal method	SnO <sub>2</sub> particles loaded on SnO nanorods	120	N/A	N/A	N/A	30
NO <sub>2</sub>	SnO <sub>2</sub> -SnO	Hydrothermal method	SnO <sub>2</sub> nanoparticles loaded on SnO nanosheets	50	19.1	100 ppm	0.1 ppm	This work

**Fig. 6** Detection selectivity of the SnO<sub>2</sub>-SnO hybrid for 100 ppm NO<sub>2</sub>, C<sub>2</sub>H<sub>5</sub>OH and CO at 50 °C. The inset shows the magnified response curves for C<sub>2</sub>H<sub>5</sub>OH and CO.**Fig. 7** Variation of the resistance ( $R_a$ ) measured in air versus time for the SnO<sub>2</sub>-SnO hybrid (top) and the SnO<sub>2</sub> (bottom), having varied temperature from approximately 20 to 200 °C (target temp.).

with rising temperature, the oxygen in air can be adsorbed on the surface of the SnO<sub>2</sub> to create chemisorbed oxygen species (O<sub>2</sub><sup>-</sup>, O<sup>-</sup>, O<sup>2-</sup>) and consume the inner electrons. Such a process could produce surface polarization energy levels and result in the increase in the bulk resistance. With the oxygen species concentration reaching saturation, the bulk resistance enters a relative stable state (bottom of Fig. 7). As shown at the top of Fig. 7, the resistance variation of the SnO<sub>2</sub>-SnO hybrid is a little more complicated than that of the pure SnO<sub>2</sub> sample. Compared with n-type SnO<sub>2</sub>, SnO belongs to the p-type semiconductors (the majority of carriers are holes).<sup>30,52</sup> With increasing temperature, the resistance of the hybrid rises first and then declines to a stable state, which can be explained as follows: in the composite, oxygen is first adsorbed on the surface of the SnO<sub>2</sub> material and the oxygen ionization process can consume large amounts of free electrons, and thus holes are left in the composites, leading to the increase in the resistance just like for the pure SnO<sub>2</sub>. With further increase in the temperature, the SnO nanocrystals become the active and main components for the chemisorption of oxygen, resulting in the further increase in hole concentration in the conduction band of the composite. The increase in hole concentration can enhance the speed of electron penetration on both sides of the p–n heterojunctions, resulting in the declining tendency of the resistance of the hybrids. Similar to that of the pure SnO<sub>2</sub> sample, the resistance reaches a relatively stable value in the end.

The n-type SnO<sub>2</sub> and p-type SnO can form a heterojunction structure at the interface, which is considered to be the primary reason for the enhanced sensing response and low operating temperature of the hybrid material.<sup>53</sup> As shown in Fig. 8A, theoretically, the forbidden band width between  $E_c$  (conduction band) and  $E_v$  (valence band) is 3.6 eV in a SnO<sub>2</sub> semiconductor and is about 2.9 eV in SnO. The p–n heterojunction leads to the formation of a depletion region. It can be seen that the bending of forbidden bands occurs at the interface zone to make the  $E_f$  (Fermi level) equilibrium between the

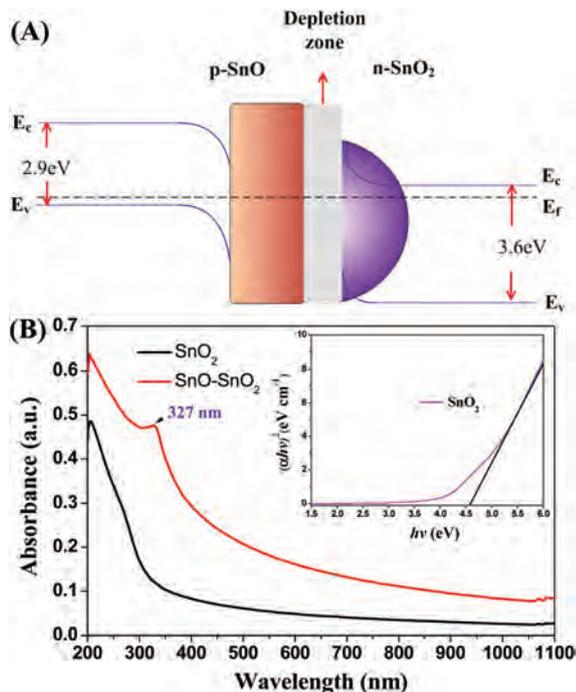


Fig. 8 (A) General schematic of the heterostructure of the n-SnO<sub>2</sub> nanoparticles on the surface of the p-SnO nanosheets with a p-n junction. (B) UV-vis spectra of the SnO<sub>2</sub>-SnO and pure SnO<sub>2</sub> materials. The inset shows the plot of  $(\alpha hv)^2$  versus  $(hv)$  from the pure SnO<sub>2</sub>.

SnO<sub>2</sub>/SnO band gaps. In general, the narrower the width of the forbidden band, the lower the required energy for the transition of electrons near the bottom of the conduction band. UV-vis spectra of the SnO<sub>2</sub>-SnO and pure SnO<sub>2</sub> materials are shown in Fig. 8B, from which the band gaps of these two samples can be evaluated. From the calculations, even though the actual forbidden band widths of these two samples are larger than that of bulk SnO<sub>2</sub> (3.6 eV) due to the quantum confinement effect, the band width of the SnO<sub>2</sub>-SnO hybrid is much lower (3.79 eV) than that of the pure SnO<sub>2</sub> (4.6 eV). The change of the band gap can well explain the reason why the operating temperature of the SnO<sub>2</sub>-SnO composite is much lower than that of the pure n-SnO<sub>2</sub>. Meanwhile, as soon as the concentration of the oxygen anions (mainly O<sup>-</sup>) gets to a saturation value, the height of the potential barrier and the width of the depletion region do not change.<sup>54</sup> The width of the depletion zone directly decides the response of a sensor at low temperature ranges. When the sensors are exposed to NO<sub>2</sub> gas (oxidizing gas), NO<sub>2</sub> molecules are absorbed on the surface of both SnO<sub>2</sub> and uncovered SnO. Similarly to the pure SnO<sub>2</sub> sensor, the absorbed NO<sub>2</sub> can capture the electrons from the conduction bands of the semiconductors, resulting in the increase in the resistance of the SnO<sub>2</sub>-SnO composite which can be seen in Fig. 5B. After trapping electrons from the conduction band and oxygen anions, NO<sub>2</sub> gas molecules became NO<sub>2</sub><sup>-</sup>. At the same time, the chemisorbed oxygen species were released in air once more. When the surface of the sample was

almost fully covered by NO<sub>2</sub><sup>-</sup>, the mobility of the free electrons on the surface was largely reduced, making the resistance of the composite increase to a saturation value. Meanwhile, the reduced mobility of the free electrons and the increased hole concentration could lead to the increase in the thickness of the depletion region, which may be another reason for the increase in sensor resistance.<sup>47</sup>

When NO<sub>2</sub> was adsorbed on the pure SnO<sub>2</sub>, owing to the strong oxidability of NO<sub>2</sub>, the electrons captured by NO<sub>2</sub> could be hardly exchanged with fresh oxygen, which may have been responsible for the poor recovery of the pure SnO<sub>2</sub>. On the other hand, when the pure SnO<sub>2</sub> sensor was exposed to NO<sub>2</sub> gas at a high temperature, because of the high electron affinity, NO<sub>3</sub><sup>-</sup> could be formed on the surface of the composites, and meanwhile, abundant holes were left in the conduction band and became the major carriers. Therefore, SnO<sub>2</sub> showed p-type semiconductor properties and the resistance showed a decrease in the process. Such an interesting phenomenon (called percolation behavior) has also been reported by Neri *et al.*<sup>51,55</sup>

## 4. Conclusions

The gas sensing properties of a SnO<sub>2</sub>-SnO composite material with a p-n heterojunction structure have been investigated for the detection of NO<sub>2</sub> gas. The results showed that the sensing performance of the SnO<sub>2</sub>-SnO hybrid is much higher than that of pure SnO<sub>2</sub> nanoparticles. The optimum temperature of the SnO<sub>2</sub>-SnO sensor is as low as 50 °C, at which the sensor showed excellent linearity between the sensing response and the concentration of NO<sub>2</sub> gas, with a wide concentration range (15 to 110 ppm), low LOD (*ca.* 0.1 ppm) and good selectivity. Even though the pure SnO<sub>2</sub> nanoparticles possessed small crystalline grain sizes and large specific surface areas, the sensing performance of the pure SnO<sub>2</sub> for NO<sub>2</sub> gas detection was still lower than that of the SnO<sub>2</sub>-SnO hybrids. In the proposed sensing mechanism, the p-n heterojunction structure formed at the interface of the SnO<sub>2</sub>-SnO may account for the improved sensing performance, especially the much lowered operating temperature. The present study shows that, in addition to optimizing the crystal size and improving the specific surface area, fabricating semiconductor materials with a p-n heterojunction structure is also a very efficient way to design gas sensors with high sensing performances.

## Acknowledgements

This work was supported by the National Natural Science Foundation of China (no. 21275136), the Natural Science Foundation of Jilin province, China (no. 201215090), Jilin province postdoctoral scientific research project (no. RB201317) and Jilin province education department scientific research project (no. 201592).

## Notes and references

- 1 H. K. Wang and A. L. Rogach, *Chem. Mater.*, 2014, **26**, 123–133.
- 2 Z. W. Chen, D. Y. Pan, Z. Li, Z. Jiao, M. H. Wu, C. H. Shek, C. M. L. Wu and J. K. L. Lai, *Chem. Rev.*, 2014, **114**, 7442–7486.
- 3 M. D'Arienzo, D. Cristofori, R. Scotti and F. Morazzoni, *Chem. Mater.*, 2013, **25**, 3675–3686.
- 4 T. Kida, S. Fujiyama, K. Suematsu, M. Yuasa and K. Shimano, *J. Phys. Chem. C*, 2013, **117**, 17574–17582.
- 5 K. Suematsu, Y. Shin, Z. Q. Hua, K. Yoshida, M. Yuasa, T. Kida and K. Shimano, *ACS Appl. Mater. Interfaces*, 2014, **6**, 5319–5326.
- 6 X. L. Wang, N. Aroonyadet, Y. Z. Zhang, M. Meckenburg, X. Fang, H. T. Chen, E. Goo and C. W. Zhou, *Nano Lett.*, 2014, **14**, 3014–3022.
- 7 N. D. Chinh, N. V. Toan, V. V. Quang, N. V. Duy, N. D. Hoa and N. V. Hieu, *Sens. Actuators, B*, 2014, **201**, 7–12.
- 8 Y. J. Chen, L. Nie, X. Y. Xue, Y. G. Wang and T. H. Wang, *Appl. Phys. Lett.*, 2006, **88**, 083105.
- 9 Y. J. Chen, X. Y. Xue, Y. G. Wang and T. H. Wang, *Appl. Phys. Lett.*, 2005, **87**, 233503.
- 10 Z. Lou, L. L. Wang, R. Wang, T. Fei and T. Zhang, *Solid-State Electron.*, 2012, **76**, 91–94.
- 11 Y. Liu, Y. Jiao, Z. Zhang, F. Qu, A. Umar and X. Wu, *ACS Appl. Mater. Interfaces*, 2014, **6**, 2174–2184.
- 12 L. W. Wang, S. R. Wang, Y. S. Wang, H. X. Zhang, Y. F. Kang and W. P. Huang, *Sens. Actuators, B*, 2013, **188**, 85–93.
- 13 L. Mei, J. W. Deng, X. M. Yin, M. Zhang, Q. H. Li, E. D. Zhang, Z. Xu, L. B. Chen and T. H. Wang, *Sens. Actuators, B*, 2012, **166**, 7–11.
- 14 C. P. Gu, X. J. Xu, J. R. Huang, W. Z. Wang, Y. F. Sun and J. H. Liu, *Sens. Actuators, B*, 2012, **174**, 31–38.
- 15 A. V. Tadeev, G. Delabouglise and M. Labeau, *Thin Solid Films*, 1999, **337**, 163–165.
- 16 P. Ivanov, E. Llobet, X. Vilanova, J. Brezmes, J. Hubalek and X. Correig, *Sens. Actuators, B*, 2004, **99**, 201–206.
- 17 G. Neri, A. Bonavita, G. Micali, N. Donato, F. A. Deorsola, P. Mossino, I. Amato and B. De Benedetti, *Sens. Actuators, B*, 2006, **117**, 196–204.
- 18 X. Y. Xue, Z. H. Chen, C. H. Ma, L. L. Xing, Y. J. Chen, Y. G. Wang and T. H. Wang, *J. Phys. Chem. C*, 2010, **114**, 3968–3972.
- 19 Y. C. Lee, H. Huang, O. K. Tan and M. S. Tse, *Sens. Actuators, B*, 2008, **132**, 239–242.
- 20 T. Zhang, L. Liu, Q. Qi, S. C. Li and G. Y. Lu, *Sens. Actuators, B*, 2009, **139**, 287–291.
- 21 P. Sun, Y. S. Yu, J. Xu, Y. F. Sun, J. Ma and G. Y. Lu, *Sens. Actuators, B*, 2011, **160**, 244–250.
- 22 B. Kim, Y. J. Lu, A. Hannon, M. Meyyappan and J. Li, *Sens. Actuators, B*, 2013, **177**, 770–775.
- 23 C. Jin, H. Kim, S. Park, H. W. Kim, S. Lee and C. Lee, *Ceram. Int.*, 2012, **38**, 6585–6590.
- 24 F.-C. Chung, R.-J. Wu and F.-C. Cheng, *Sens. Actuators, B*, 2014, **190**, 1–7.
- 25 J. Zhang, X. H. Liu, S. H. Wu, M. J. Xu, X. Z. Guo and S. R. Wang, *J. Mater. Chem.*, 2010, **20**, 6453–6459.
- 26 S. W. Choi, A. Katoch, G. J. Sun, P. Wu and S. S. Kim, *J. Mater. Chem. C*, 2013, **1**, 2834–2841.
- 27 R.-J. Wu, D.-J. Lin, M.-R. Yu, M. H. Chen and H.-F. Lai, *Sens. Actuators, B*, 2013, **178**, 185–191.
- 28 M. M. Liu, R. Z. Zhang and W. Chen, *Chem. Rev.*, 2014, **114**, 5117–5160.
- 29 F. F. Ren, C. Q. Wang, C. Y. Zhai, F. X. Jiang, R. R. Yue, Y. K. Du, P. Yang and J. K. Xu, *J. Mater. Chem. A*, 2013, **1**, 7255–7261.
- 30 A. Shanmugasundaram, P. Basak, L. Satyanarayana and S. V. Manorama, *Sens. Actuators, B*, 2013, **185**, 265–273.
- 31 S. W. Choi, A. Katoch, J. Zhang and S. S. Kim, *Sens. Actuators, B*, 2013, **176**, 585–591.
- 32 M. Mashock, K. H. Yu, S. M. Cui, S. Mao, G. H. Lu and J. H. Chen, *ACS Appl. Mater. Interfaces*, 2012, **4**, 4192–4199.
- 33 N. D. Khoang, D. D. Trung, N. V. Duy, N. D. Hoa and N. V. Hieu, *Sens. Actuators, B*, 2012, **174**, 594–601.
- 34 F. Shao, M. W. G. Hoffmann, J. D. Prades, R. Zamani, J. Arbiol, J. R. Morante, E. Varechikina, M. Rumyantseva, A. Gaskov, I. Giebelhaus, T. Fischer, S. Mathur and F. Hernandez-Ramirez, *Sens. Actuators, B*, 2013, **181**, 130–135.
- 35 Z. J. Wang, Z. Y. Li, J. H. Sun, H. N. Zhang, W. Wang, W. Zheng and C. Wang, *J. Phys. Chem. C*, 2010, **114**, 6100–6105.
- 36 T. Kida, A. Nishiyama, Z. Hua, K. Suematsu, M. Yuasa and K. Shimano, *Langmuir*, 2014, **30**, 2571–2579.
- 37 P. Sun, Y. X. Cai, S. S. Du, X. M. Xu, L. You, J. Ma, F. M. Liu, X. S. Liang, Y. F. Sun and G. Y. Lu, *Sens. Actuators, B*, 2013, **182**, 336–343.
- 38 W. W. Guo, T. M. Liu, R. Sun, Y. Chen, W. Zeng and Z. C. Wang, *Mater. Lett.*, 2012, **89**, 5–8.
- 39 G. W. She, X. Huang, L. L. Jin, X. P. Qi, L. X. Mu and W. S. Shi, *Small*, 2014, **10**, 4685–4692.
- 40 K. T. Lee, C. H. Lin and S. Y. Lu, *J. Phys. Chem. C*, 2014, **118**, 14457–14463.
- 41 Y. Chen, J. Lu, S. Wen, L. Lu and J. M. Xue, *J. Mater. Chem. A*, 2014, **2**, 17857–17866.
- 42 S. H. Ahn, D. J. Kim, W. S. Chi and J. H. Kim, *Adv. Funct. Mater.*, 2014, **24**, 5037–5044.
- 43 L. Li, S. J. He, M. M. Liu, C. M. Zhang and W. Chen, *Anal. Chem.*, 2015, **87**, 1638–1645.
- 44 H. Zhang, J. C. Feng, T. Fei, S. Liu and T. Zhang, *Sens. Actuators, B*, 2014, **190**, 472–478.
- 45 A. Sharma, M. Tomar and V. Gupta, *Sens. Actuators, B*, 2013, **176**, 875–883.
- 46 H. J. Song, L. C. Zhang, C. L. He, Y. Qu, Y. F. Tian and Y. Lv, *J. Mater. Chem.*, 2011, **21**, 5972–5977.
- 47 F. Song, H. L. Su, J. J. Chen, W. J. Moon, W. M. Lau and D. Zhang, *J. Mater. Chem.*, 2012, **22**, 1121–1126.
- 48 H. L. Huang, Y. J. Liu, Q. Z. Gao, W. S. Ruan, X. M. Lin and X. Li, *ACS Appl. Mater. Interfaces*, 2014, **6**, 10258–10264.

- 49 C. N. Xu, J. Tamaki, N. Miura and N. Yamazoe, *Sens. Actuators, B*, 1991, **3**, 147–155.
- 50 R. K. Paul, S. Badhulika, N. M. Saucedo and A. Mulchandani, *Anal. Chem.*, 2012, **84**, 8171–8178.
- 51 G. Neri, S. G. Leonardi, M. Latino, N. Donato, S. Baek, D. E. Conte, P. A. Russo and N. Pinna, *Sens. Actuators, B*, 2013, **179**, 61–68.
- 52 A. K. Sinha, P. K. Manna, M. Pradhan, C. Mondal, S. M. Yusuf and T. Pal, *RSC Adv.*, 2014, **4**, 208–211.
- 53 X. H. Zhou, Q. X. Cao, H. Huang, P. Yang and Y. Hu, *Mater. Sci. Eng., B*, 2003, **99**, 44–47.
- 54 L. I. Trakhtenberg, G. N. Gerasimov, V. F. Gromov, T. V. Belysheva and O. J. Ilegbusi, *Sens. Actuators B*, 2012, **169**, 32–38.
- 55 N. Savage, B. Chwieroth, A. Ginwalla, B. R. Patton, S. A. Akbar and P. K. Dutta, *Sens. Actuators B*, 2001, **79**, 17–27.
- 56 S. F. Zhang, F. Ren, W. Wu, J. Zhou, X. H. Xiao, L. L. Sun, Y. Liu and C. Z. Jiang, *Phys. Chem. Chem. Phys.*, 2013, **15**, 8228–8236.
- 57 D. Hu, B. Q. Han, S. J. Deng, Z. P. Feng, Y. Wang, J. Popovic, M. Nuskol, Y. D. Wang and I. Djerdj, *J. Phys. Chem. C*, 2014, **118**, 9832–9840.